

# Effect of the Magnus Force in the PITCHf/x Tracking System

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A new method is proposed for calculating the  $pfx_x$  and  $pfx_z$  parameters in the PITCHf/x tracking system.

## I. INTRODUCTION

The PITCHf/x system uses two cameras to track pitches between pitcher and batter, determining the coordinates of the ball at 1/60-sec intervals. The resulting trajectory  $x(t), y(t), z(t)$  ( $t$  is the time) is fit to a nine-parameter (or “9P”) fit corresponding to constant acceleration in each of the three dimensions. All quantities reported in the PITCHf/x data base, such as the pitch speed, the location of the pitch as it crosses the plate, the “break” of the pitch, etc., are derived from the fitted trajectory rather than from the original data. The nine parameters are the three initial positions  $x_0$ ,  $y_0$ , and  $z_0$ ; the three initial velocities  $v_{x0}$ ,  $v_{y0}$ ,  $v_{z0}$ ; and the three accelerations  $a_x$ ,  $a_y$ , and  $a_z$ . Here the coordinates refer to the usual PITCHf/x coordinate system, where the origin is at the point of home plate,  $\hat{y}$  points towards the pitcher,  $\hat{z}$  points vertically upward, and  $\hat{x} = \hat{y} \times \hat{z}$  (i.e., the  $x$  axis points to the catcher’s right). The 9P fit is an approximation to the actual equations of motion,

$$\begin{aligned}\ddot{x} &= -KC_D vv_x - KC_L vv_y \sin \phi \\ \ddot{y} &= -KC_D vv_y + KC_L v (v_x \sin \phi - v_z \cos \phi) \\ \ddot{z} &= -KC_D vv_z + KC_L vv_y \cos \phi - g.\end{aligned}\tag{1}$$

Here  $g$  is the acceleration due to gravity (32.174 ft/s<sup>2</sup>),  $C_D$  and  $C_L$  are the drag and lift coefficients, respectively, and  $K = 5.44 \times 10^{-3}$  ft<sup>-1</sup> is a numerical factor.[1] In these expressions, the spin axis is assumed to lie in the  $x - z$  plane and makes an angle  $\phi$  with the  $x$  axis, with a sign such that  $\phi = 90^\circ$  corresponds to the spin pointing upward, along the  $z$  axis. Noting that  $v_y$  is negative, it is easy to see that the Magnus force makes an angle  $\theta = \phi - 90^\circ$  with the  $x$  axis. Therefore  $\phi = 0^\circ$  (topspin) results in a downward acceleration, and  $\phi = 90^\circ$  (sidespin) results in an acceleration to the catcher’s right, exactly as expected.

Two quantities calculated by PITCHf/x are  $pfx_x$  and  $pfx_z$ , the deviation of the pitch trajectory in the  $x$  and  $z$  directions from that expected in the absence of the Magnus force, as measured between  $y=40$  ft and the front edge of home plate,  $y=1.417$  ft. The exact way to calculate these quantities is to compare the actual trajectory with that computed by solving the equations of motion, Eqs. 1, with  $C_L=0$ . The PITCHf/x system uses an approximate method given by the prescription

$$\begin{aligned}pfx_x &= \frac{1}{2} a_{xM} t_{40}^2 \\ pfx_z &= \frac{1}{2} a_{zM} t_{40}^2,\end{aligned}\tag{2}$$

where  $t_{40}$  is the time of flight between  $y = 40$  and  $y = 1.417$  ft and the accelerations due to the Magnus force,  $a_{xM}$  and  $a_{zM}$ , are given by

$$\begin{aligned}a_{xM} &= a_x \\ a_{zM} &= a_z + g.\end{aligned}\tag{3}$$

This prescription assumes that  $a_x$  and  $a_z + g$  are due entirely to the Magnus force. An inspection of Eq. 1 shows that this is a good approximation to the extent that the drag-related terms (i.e., those proportional

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to  $C_D$ ) can be neglected in the expressions for the  $x$  and  $z$  accelerations. In this brief note, an alternate approximate prescription for  $pfx_x$  and  $pfx_z$  is proposed, then compared to the earlier prescription as well as to an exact calculation for a large selection of actual pitches.

## II. ALTERNATE TECHNIQUE

The alternate technique uses an expression identical to Eq. 2 but with accelerations  $a_{xM}$  and  $a_{zM}$  modified to approximately remove the effects of drag as follows:

$$\begin{aligned} a_{xM} &= a_x - a_y \frac{\langle v_x \rangle}{\langle v_y \rangle} \\ a_{zM} &= a_z - a_y \frac{\langle v_z \rangle}{\langle v_y \rangle} + g, \end{aligned} \quad (4)$$

where the brackets indicate the time average over the full trajectory. The rationale behind this improved approximation is found in Eq. 1, where it is observed that the contribution of the drag to each component of the instantaneous acceleration is directly proportional to that component of the instantaneous velocity. The  $y$  component of the drag acceleration is approximated by  $a_y$ , an approximation that is expected to be quite good since  $|v_x/v_y| \ll 1$  and  $|v_z/v_y| \ll 1$ . To get the time-averaged drag, the time-averaged velocities are needed and calculated using straightforward kinematics. First  $v_{yf}$ , the  $y$  component of velocity at  $y_f = 1.417$  ft is calculated. Then the total flight time  $T$  is calculated. Then  $a_x$  and  $a_z$  are used to calculate  $v_{xf}$  and  $v_{zf}$ . Finally, the initial and final velocities are used to calculate the time-averaged values. The particular sequence of equations is as follows:

$$\begin{aligned} v_{yf} &= -\sqrt{v_{y0}^2 + 2a_y(y_f - y_0)} \\ T &= \frac{v_{yf} - v_{y0}}{a_y} \\ v_{xf} &= v_{x0} + a_x T \\ v_{zf} &= v_{z0} + a_z T, \end{aligned} \quad (5)$$

and

$$\langle v_x \rangle = (v_{xf} + v_{x0})/2 \quad (6)$$

and similarly for  $\langle v_y \rangle$  and  $\langle v_z \rangle$ . As an aside, it is noted that the angle  $\phi$  can be found approximately from

$$\phi = \arctan\left(\frac{a_{zM}}{a_{xM}}\right) + 90^\circ. \quad (7)$$

## III. COMPARING THE TWO TECHNIQUES

To facilitate the notation, the quantities  $\Delta x$  and  $\Delta z$  are defined to be the difference between the new and old values of  $pfx_x$  and  $pfx_z$ , respectively. Also  $pfx$  is defined as  $\sqrt{pfx_x^2 + pfx_z^2}$ . The results of the analysis are shown in Fig. 1, where nearly 8000 pitches from games played in Toronto during the period April-June, 2007 are analyzed. The  $\Delta z$  plot shows a consistent systematic difference between the two techniques which is easily understood. Since the ball always follows a downward trajectory, the  $z$  component of drag is always upward (positive). Using the old system of calculating  $pfx_z$ , the upward drag produces a deflection which is systematically more positive than the exact value. That is,  $\Delta z$  (new value minus old value) is systematically negative. The effect on the  $x$  coordinate depends on the direction of the  $x$  velocity, and the two plots show a systematic shift of  $pfx_x$  in the positive direction (for a negative  $v_{x0}$ ) or the negative direction (for a positive  $v_{x0}$ ), with the reasoning being exactly the same as for the  $z$  deflection. The final plot shows the correlation between the total deflection  $pfx$  and the inferred value of the lift coefficient  $C_L$  for these trajectories, the latter calculated by doing a non-linear least-squares fit to the smoothed trajectories using the full equations

of motion. Note that  $pfx$  should be linearly proportional to  $C_L$  and independent of the initial velocity. This latter point can be seen from Eq. 2, since the Magnus acceleration is proportional to  $v^2$  but the time is proportional to  $1/v$ . The figure shows that the values  $pfx$  calculated with Eq. 4 are perfectly correlated with  $C_L$ , indicating the method is a very good approximation to the exact solution. Indeed, the red curve is a nearly perfect straight line passing through the origin. On the other hand, the values calculated using Eq. 3 are less well correlated, as indicated by the scatter in the values.

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[1] The constant  $K$  scales with air density. The value given assumes normal temperature and pressure.

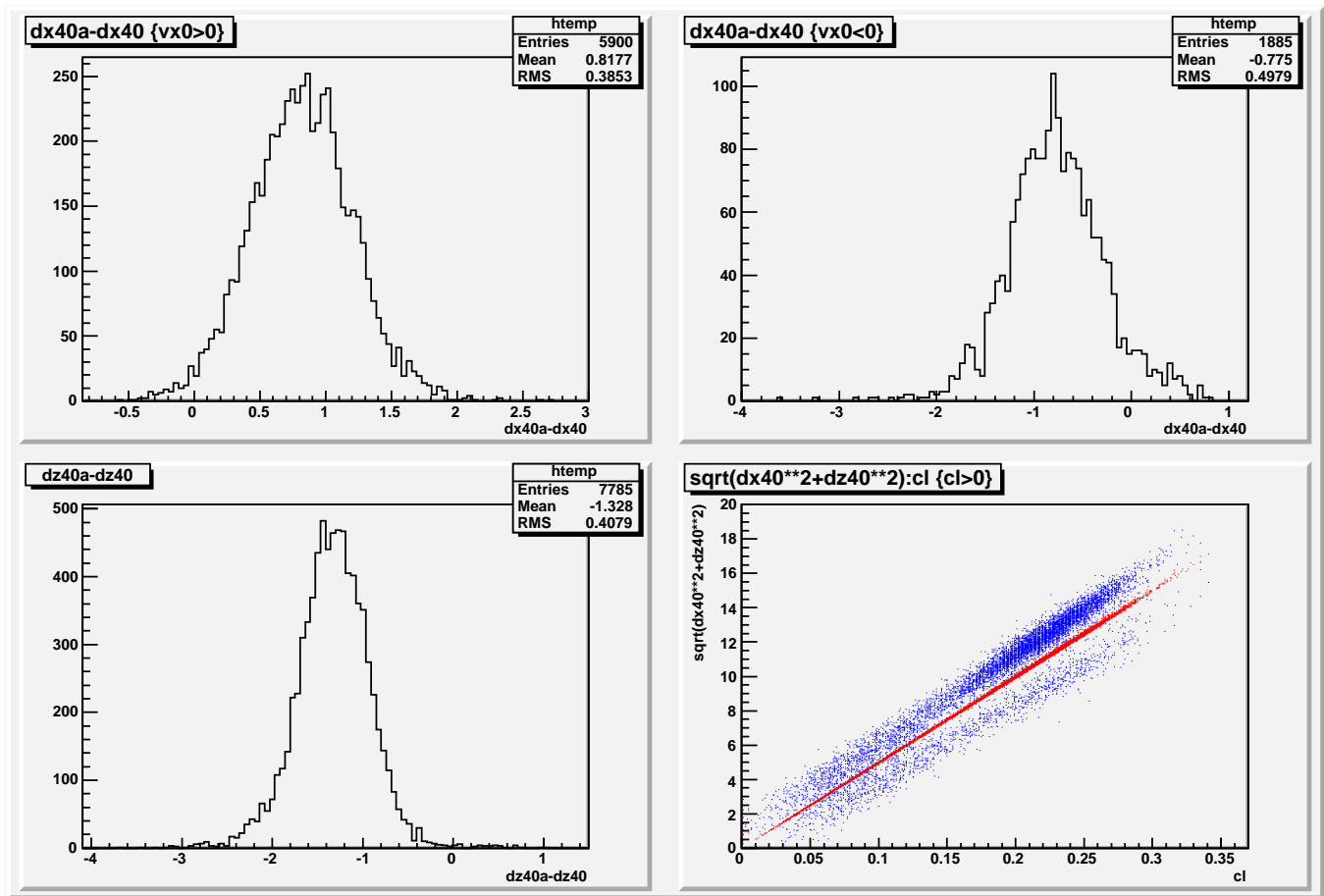


FIG. 1: (left) Plots of 7785 pitches from the Toronto PITCH/x data. The upper row is a plot of  $\Delta x$  for  $v_{x0} > 0$  (left) and  $v_{x0} < 0$  (right). The lower left plot is  $\Delta z$ . The lower right plot shows the correlation between  $pf_x$  and  $C_L$ , with the blue points and red points calculated using the prescriptions of Eq. 3 and Eq. 4, respectively.